Lecture 20

Thus
$$I_p = I_0 \sum_{j,j'} e^{-i(\vec{k} - \vec{k}_i) \cdot (\vec{r}_j - \vec{r}_{j'})}$$

(i) Forward scattering : $\vec{k} = \vec{k}_i$

$$I_p = I_0 \left(\sum_{i=1}^{N} 1\right) \left(\sum_{i'=1}^{N} 1\right) = N^2 I_0$$

Just as we saw when considering small ($\ll \lambda$) collections of scatterers, <u>coherent</u> scattering goes like N^2

(ii) Side scattering: $\vec{k} \neq \vec{k_i}$

Define $\vec{\Delta k} = \vec{k} - \vec{k_i}$ as the "scattering wavevector"

We've been assuming elastic scattering (i.e. no absorption), so that

We'll consider 2 cases-dipole scatterers arranged randomly or in a perfect crystal(a) <u>Random</u> scatters

 \Rightarrow For any atom j, the sum

$$\sum_{j'=1}^{N} e^{-i\overrightarrow{\Delta k} \cdot (\vec{r}_{j} - \vec{r}_{j'})} = e^{-i\overrightarrow{\Delta k} \cdot \vec{r}_{j}} \underbrace{\sum_{j'=j}^{N} e^{i\overrightarrow{\Delta k} \cdot \vec{r}_{j'}}}_{\sum_{j'=j} + 1} + 1$$

This term oscillates wildly and

randomly as the sum over j'

is made, and thus averages

to zero
$$(\int_{0}^{2\pi} e^{i\varphi} d\varphi = 0)$$

Thus $I_p = I_0 \sum_{i=1}^{N} = NI_0$

Thus we recover the case of N "<u>independent</u>" scatters when they are <u>randomly positioned</u>. We see that incoherent scattering goes like $I = NI_0$

For randomly placed molecules/scatters, we do have side scattering and thus attenuation according to the exponential law

$$I(z) = I_{inc} e^{-\sigma N_z}$$

Where N is the number density and $\,\sigma\,$ is the single-molecule scattering cross section

(b) Scattering centers arranged on a perfect lattice

This case can be shown to be equivalent to the perfectly homogeneous medium case: When $\lambda \gg$ atomic spacing, only coherent formed scattering occurs.

Real life: <u>Thermal fluctuations (defects)</u> => imperfect lattice, in addition to the <u>finite size of</u> the material (or the incident beam)

 \Rightarrow Recover exponential attention

Wavelength dependence of scattered light

We must go back to the form of the cross section

$$\sigma(\omega) = \frac{\omega^4 \alpha^2(\omega)}{12\pi \varepsilon_0^2 c^4}$$

Clearly, if $\alpha(\omega)$ shows strong resonances, as it does in the CEO model where

$$\alpha_{CEO}(\omega) = \frac{e^2 / m}{\omega_0^2 - \omega^2 + i\omega\gamma}$$

Then the scattering will also be strongly resonant. (e.g. color of metal nanoparticles)

Rayleigh scattering

This situation typically involves <u>nonresonant scattering</u> in <u>transparent media</u>. A classical example is the <u>propagation</u> of <u>light</u> in <u>air</u>. The various molecules in the atmosphere have their absorption resonances in the ultraviolet. We thus have the situation for visible light at ω :

$$\omega \ll \omega_0, \omega \gamma \ll \omega_0^2$$

This gives

$$\alpha_{CEO}(\omega) \simeq \frac{e^2}{m\omega_0^2}$$

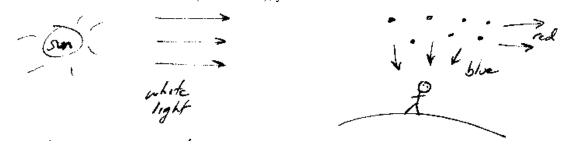
Which is independent of frequency.

Thus
$$\sigma(\omega) = \frac{e^2 \omega^4}{12\pi \varepsilon_0^2 m \omega_0^2 c^4}$$

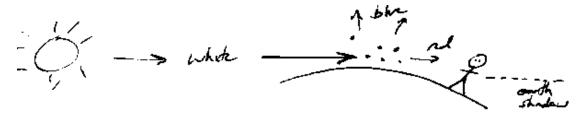
If the polarizability of the scattering molecules is not too strong (as is the case for visible light in transparent media), then the <u>frequency dependence of light scattering is</u> ω^4 .

Note that scattering of blue (near UV) light at 400 nm is <u>sixteen times</u> as strong as scattering of red/near-IR light at 800 nm. This gives us the classical argument for the <u>blue sky</u>, similar to that first developed by Lord Rayleigh:

- (1) The frequency dependence of the index of refraction is fairly weak across the visible spectrum
- (2) The molecules in the atmosphere are completely randomly positioned, and thus show density fluctuations on the scale of a wavelength. These density fluctuations cause light to scatter on propagation through the atmosphere
- (3) The ω^4 law means blue is scattered more than red. Thus, when looking in a direction different from that of the source (i.e. the sum), you see blue!



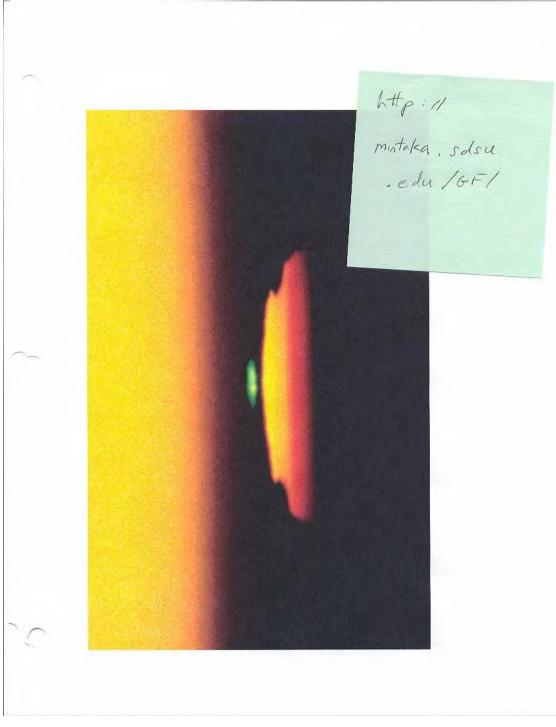
Note that this also explains the red sunset :



As the light travel through a long distance in the atmosphere at sunset, most of the blue is scattered away ,leaving only the red. (If you look carefully after look on a clear evening,you will notice a continuous shift in the color of the sky from red to blue as you look west to east.) note Mount Pinatubo effect !

So why are clouds white (when you are looking at reflected light)?? Clouds are collections of water droplets, which have a size $\geq \lambda$ for visible light .Therefore all visible wavelengths are scattered with roughly equal efficiency, and thus you see white.

Why is the sky gray on a cloudy day? If the sun is behind clouds, you are looking at transmitted light, which is attenuated by scattering. The attenuation is independent of λ , and attenuated white is just gray!



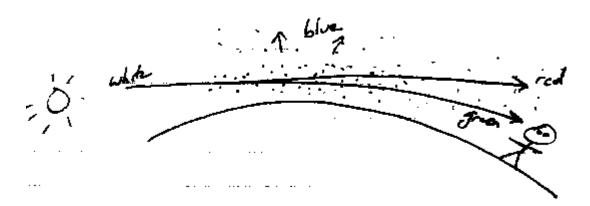
Back to sunsets: at the very moment just after the top of the sun dips below the horizon ,the "<u>green flash</u> "may sometimes be observed .It can only be observed when the sun sets over the ocean or a very flat plain (sometimes from airplanes ,too).

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The green flash is caused by a combination of effects:

- (1) At sunset, the final rays you see travel a long way through the atmosphere => the blue is scattered away.
- (2) The yellow and orange are weakly absorbed by water vapor and possibly oxygen molecules; over a long distance this absorption may be significant \rightarrow only red+green left.
- (3) As we have seen, rays in the atmosphere can bend, due to the inhomogeneous (radial ,in this case)

Index of refraction. Due to dispersion, the red bends less than the green:



Another practical consequence of Rayleigh scattering: intrinsic losses in optical fibers.

Optical fibers are generally made of glass (e.g. silica), which is a disordered medium. In fact, it is an "amorphous" solid, with the same sort of density fluctuations you find in a gas. These density fluctuations cause light scattering out of the fiber, and since they are unavoidable, they set the lower limit to possible fiber attenuation.

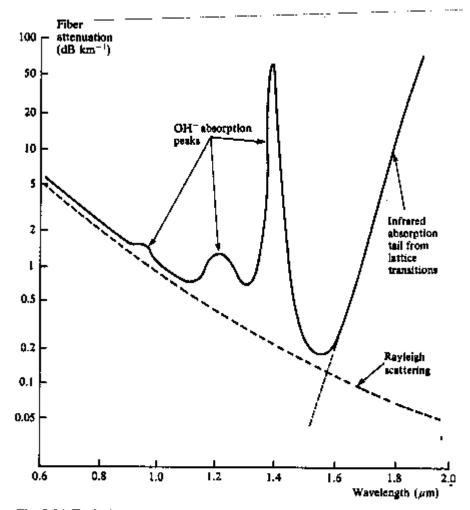
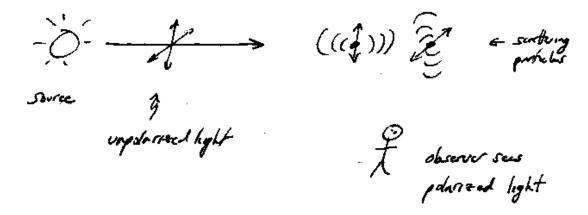


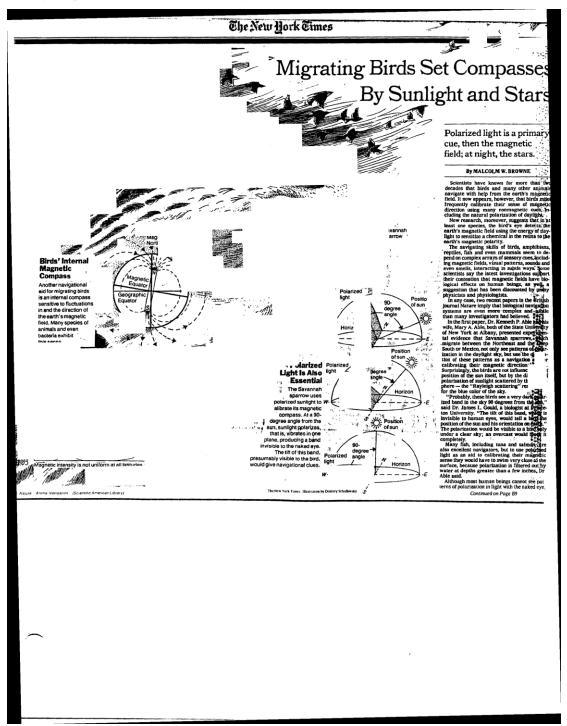
Fig. 8.24 Typical attenuation versus wavelength plot for a silica-based optical fiber. The contribution from Rayleigh scattering is shown, as are the other two main loss mechanisms, namely the infrared absorption tail and the hydroxyl (OH^-) absorption peaks.

Polarization of scattered light

- A consequence of the dipole na## of the scattering



(you can check this by looking at the sky through Polaroid sunglasses and rotating them. You should look at an angle $\sim 90^{\circ}$ from the sun.)



Is this of any use? It is if you are a bird! It has long been know that bird navigate with the help of the earth's magnetic field. Now it appears that some birds "calibrate" their magnetic field sense by sensing the polarization of scattered sunlight. (see NY Times 9/29/93).

Reading : Guenther chap.5(except resonators + guided waves)